

1898
Problem NO 1

1. Polarization of conducting sphere.

Consider the Laplace Equation

$$\nabla^2 \varphi = 0$$

And we have the boundary condition

$$r \rightarrow \infty, \varphi \rightarrow -Er \cos \theta ; \quad r \rightarrow R \quad \frac{\partial \varphi}{\partial r} \rightarrow 0$$

According to the Legendre polynomials, we have the form of the solution

$$\varphi = -Er \cos \theta + \frac{A}{r^2} \cos \theta \quad A = ER^3$$

$$\text{i.e. } \varphi = -Er \cos \theta + \frac{ER^3}{r^2} \cos \theta$$

The second part is similar to the electric field generated by electric dipoles.

So, the ball is equivalent to the electric dipole

$$\vec{p} = 4\pi \epsilon_0 R^3 \vec{E}$$

2. The motion of one ball.

Since $L \gg R$, and the interaction force between

two balls $F \propto \frac{p^2}{L^4} \propto \frac{R^6}{L^4}$, it's necessary to

consider the first order induction between the balls.

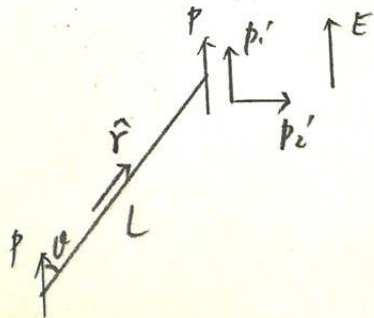
As the electric field generated by electric dipoles is,

$$\varphi_1 = \frac{\vec{p} \cdot \hat{r}}{4\pi \epsilon_0 r^2}$$

We can obtain that.

$$\vec{E}_1 = -\nabla \varphi_1 = \frac{3(\vec{p} \cdot \hat{r})\hat{r} - \vec{p}}{4\pi \epsilon_0 r^3}$$

$$\text{So, } p_1 = \frac{R^3}{L^2} p (3 \cos^2 \theta - 1) \quad p_2' = \frac{R^3}{L^2} p \cdot 3 \cos \theta \sin \theta$$



Consider the part from the center of mass to one ball.

$\vec{M}_p^{(e)}$: external electric field to the dipole

\vec{M}_{pp} : dipole to dipole

\vec{F}_{pp} : the force of dipole to dipole

$\vec{M}_p^{(i)}$: external electric field to the induction dipole.

It's clear that the total torque is

$$\vec{M}_t = \vec{M}_p^{(e)} + \vec{M}_{pp} + \left(\frac{L}{2} \hat{r}\right) \times \vec{F}_{pp} + \vec{M}_p^{(i)}$$

$$\vec{M}_p^{(e)} = \vec{p} \times \vec{E} = 0$$

$$\vec{M}_{pp} = \vec{p} \times \vec{E}_1 = \frac{3p^2 \sin\theta \cos\theta}{4\pi\epsilon_0 L^3} \hat{z} \quad (\hat{z} = \hat{r} \times \hat{\theta})$$

$$\begin{aligned} \vec{F}_{pp} &= \vec{p} \times \nabla \vec{E}_1 = \frac{3p^2}{4\pi\epsilon_0 L^4} [(\hat{p} \cdot \hat{r}) \hat{p} + (\hat{p} \cdot \hat{r}) \hat{p} + (\hat{p} \cdot \hat{p}) \hat{r} - 5 (\hat{p} \cdot \hat{r}) (\hat{p} \cdot \hat{r}) \hat{r}] \\ &= \frac{3p^2}{4\pi\epsilon_0 L^4} (2\cos\theta \hat{p} + \hat{r} - 5\cos^2\theta \hat{r}) \end{aligned}$$

$$\vec{M}'_{pp} = \frac{L}{2} \hat{r} \times \vec{F}_{pp} = - \frac{3p^2 \sin\theta \cos\theta}{4\pi\epsilon_0 L^3} \hat{z}$$

$$\vec{M}_p^{(i)} = (\vec{p}_1 + \vec{p}_2) \times \vec{E} = \vec{p}_2 \times \vec{E} = - \frac{R^3}{L^3} p E \cdot 2 \cos\theta \sin\theta \hat{z}$$

So, we can obtain the equation from angular momentum theorem.

$$M \cdot \left(\frac{L}{2}\right)^2 \ddot{\theta} = - \frac{3R^3}{L^3} p E \cos\theta \sin\theta$$

when $\theta \ll 1$, it turns to

$$M \left(\frac{L}{2}\right)^2 \ddot{\theta} = - \frac{3R^3}{L^3} p E \theta$$

$$\text{So, } T = 2\pi \sqrt{\frac{M(L/2)^2}{3R^3 p E / L^3}} = 2\pi \sqrt{\frac{M L^5}{48\pi\epsilon_0 E R^3}} = \sqrt{\frac{\pi M L^5}{12\epsilon_0 E R^3}}$$

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